Manuel Moreno-Ibáñez, Josep M. Trigo-Rodríguez & Maria Gritsevich

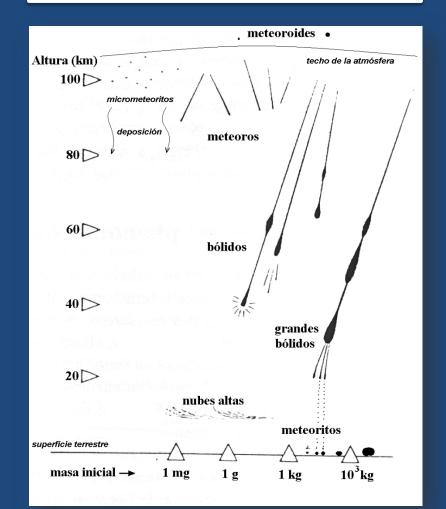


Summary

- Background
- Equations of motion
- Simplifications
- Database
- Results
- Discussion
- Conclusions
- References

- Asteroid: diameter > 10 m.
- Meteoroid: diameter < 10 m.
- Meteor: meteoroid that interacts with the Earth atmosphere.
- Fireball: meteor that is able to get deep in the atmosphere.
 Brightness similar to Venus.
- Great Fireball: bigger fireball that is able to get to lower altitudes and can reach a brightness similar to the Moon.
- Meteorites: meteor that survives to its atmospheric flight and reaches the ground.
- Micrometeorites: small grains that get the atmosphere at low velocities and are deposited on the ground.

Fireballs and Meteorites



Adapted from Rendtel et al. (1995)

Ceplecha y McCrosky (1976)

- Study of the Prairie Network to find similar meteorites to the Lost City meteorite.
- They used four criteria: End height agrees with the single-body theoretical value, calculated using *dynamic mass*, as well as with that of Lost City too within ±1.5km, when scaled for mass, velocity and entry angle in accordance with classic meteor theory.

Stulov et al. (1995), Stulov (1997) and Gritsevich (2007)

- Study of the Prairie Network to distinguish between ordinary and carbonaceous chondrites.
- The end height as the principal discrimating observational parameter in their discussion (*photometric mass*).
- Empirical Criterium:

 $PE = \log \rho_E + A \log m_{\infty} + B \log V_{\infty} + C \log(\cos Z_R)$

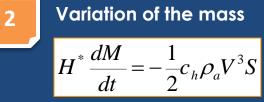
• A, B and C obtained by Least–Squared fit on PN data.

Wetherill and Revelle (1981)

- Instead of using the average values as input parameters, they gathered all the unknowns into two new variables, α and β (ballistic coefficient and mass loss parameter).
- Adjusting the equation to the registered values these new variables can be obtained.
- This describes in detail the meteoroid trajectory and allow to invent a classification for possible impacts.

Equations of motion

The equations of motion for a meteoroid entering the atmosphere projected onto the tangent and to the normal to the trajectory



$$M \frac{dV}{dt} = -\frac{1}{2} c_d \rho_a V^2 S + P \sin \gamma$$

$$MV\frac{d\gamma}{dt} = P\cos\gamma - \frac{MV^2}{R}\cos\gamma - \frac{1}{2}c_L\rho_a V^2 S$$
$$\frac{dh}{dt} = -V\sin\gamma$$

3

Extra equations

Isotermal atmosphere $\rightarrow \rho = \exp(h/h_0)$ Levin (1956) $\rightarrow S/S_e = (M/M_e)^{\mu}$



Use of dimensionless parameters $M = M_e m; V = V_e v; S = S_e s; h = h_0 y; \rho_a = \rho_0 \rho$

Where index e indicates values at the entry of the atmosphere. h0 is the scale height (7.16 km) Lifting force

Analytical solution of system: Initial conditions : $y = \infty; v = 1; m = 1$

$$m = \exp\left(-\frac{\beta(1-v^2)}{1-\mu}\right) \quad [1]$$

$$y = \ln 2\alpha + \beta - \ln \Delta, \quad \Delta = \overline{Ei}(\beta) - \overline{Ei}(\beta v^2) \quad [2]$$

$$\overline{Ei}(x) = \int_{-\infty}^{x} \frac{e^z dz}{z}$$

In this methodology we gather all the unknown values of the meteoroid's atmosphere flight motion equations into two new variables (Gritsevich, 2009):

Ballistic Coefficient

$$\alpha = \frac{1}{2}c_d \frac{\rho_0 h_0 S_e}{M_e \sin \gamma}$$

Mass loss parameter

$$\beta = (1 - \mu) \frac{c_h V_e^2}{2c_d H^*}$$

Simplifications

For quick meteors, a strong evaporation process takes place so β becomes high (β >> 1), the deceleration can be neglected and the velocity thus assumed constant. Stulov (1998, 2004) developed the following asymptotic solution:



$$v = 1$$
, $m^{1-\mu} = 1 - 2\alpha\beta e^{-y}$, $\ln 2\alpha\beta < y < \infty$ [3]

However, the meteor velocity begins to decrease in a certain vicinity of m=0. In order to account for this change in velocity we combine the Eq.[1] (valid for arbitrary β values) with the Eq. [3] suitable for high β values:

$$v = \left(\frac{\ln\left(1 - 2\alpha\beta e^{-y}\right)}{\beta} + 1\right)^{\frac{1}{2}}, \quad \ln 2\alpha\beta < y < \infty \quad [4]$$

Database

In order to test this methodology, we compare our derived terminal heights results with the fireball terminal heights registered by the Meteorite Observation and Recovery Project operated in Canada between 1970 and 1985 (MORP) (Halliday et al. 1996). We use previous α and β values derived by Gritsevich (2009).

		Table 1b.								
(1996626539153966) The results of processing of observational data for the Canadian Network fireballs.										
	188	Fireball No.	Ve ^a , km/s	α	β	$10^2\sigma$	sin γ	<i>M</i> ₁ , g	<i>M</i> ₂ , g	<i>M_{ph}</i> ^a , g
TABLE 3. Physical and orbital data for a sample of 213 fireballs. MORP ID Y M D Δ RA DE ZA ZC LH LG VI VE VH VG HB HE DU a e i q q'ω Ω S DS Q MG L MI MT DN		18	18.5	24.13	1.475	0.862	0.572	5417.24	11466.73	5600
132 2330.867 74 10 10 10 113.7 79.5 35.9 36.3 95.9 57.2 48.8 48.8 40.1 47.5 96.3 73.3 0.6 5.320 0.815 84.2 0.985 966 194.1 196.3 AI 1 9.9 2.88 29 132 2350.857 54 11 05 0 17.9 20.7 65.9 76.3 163.3 12.1.1 16.9 9.5 38.3 12.9 72.2 37.0 56 2.742 0.669 8.8 0.907 4.58 38.0 42.2 BK 2 8.7 3.83 250 93		123	16.3	37.22	1.111	0.836	0.551	1651.41	3495.55	680
141 2360739 74 11 09 R 554 220 377 386 1295 787 32.1 290 380 300 886 651. 09 2.580 0.870 2.8 0.335 4.82 2950 2260 NT AF 1 62 2.26 17 0.2 142 302 801 74 11 11 20 657 401 169 185 1474 980 225 135 393 214 749 429 15 3550 9856 165 0.70 6.38 2473 2282 E 1 7.1 2.83 140 0.8		138	16.9	38.90	2.889	2.023	0.404	3669.82	7767.96	2400
167 2472.741 75 03 01 0 1583 143 376 398 1463 942 240 240 380 212 895 253 18 2.566 0646 3.2 0.748 4.49 2592 339.7 9 DG I 10.7 4.53 3300 169 2483547 37 012 0 1550 - 146 698 750 1482 967 229 93 380 202 789 340 64 2.620 0.726 15.7 0.729 4.50 687 170.8 CM I 10.6 4.73 12000 62		141	32.1	707.19	6.403	1.243	0.767	0.09	0.19	17
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$		144	23.5	137.13	1.301	0.471	0.953	6.38	13.51	140
		169	22.9	50.52	1.575	0.601	0.337	2886.41	6109.70	12,000
		172	12.5	13.13	3.379	4.325	0.969	6916.28	14639.77	1200
		177	15.8	77.63	1.777	1.424	0.631	121.19	256.52	300
		187	18.4	83.76	3.815	2.254	0.663	83.17	176.06	340
		189	14.5	34.47	0.757	0.720	0.388	5954.14	12603.20	8100
	. Hall	192	21.0	94.01	2.458	1.115	0.860	26.95	57.05	120
251 2995824 76 98 05 R 379 543 407 409 650 380 607 607 409 959 973 821 03 11,430 0916 1155 0545 2150 1512 1326 P A 1 84 256 9 19 151 2017 77 76 14 1 84 3 584 494 497 640 410 640 660 413 588 882 758 03 1935 1095 1113 0956 3774 1545 1393 P B1 2 86 256 9	iday e	195	25.2	35.22	1.486	0.468	0.593	1563.53	3309.54	7000
265 3102022 76 11 19 0 1492 172 324 325 87 50 710 710 401 70 11022 815 03 4750 0783 1722 0985 852 1872 2370 L E 1 104 334 38 27 210 21 21 20 1153 310 256 265 1141 61 256 324 323 246 889 497 12 12 126 0944 231 0122 44 3277 2622 G EG 1 100 350 290	t al.	204	13.0	10.13	3.620	4.284	0.537	88488.54	187304.65	2900
278 3142.051 76 12 29 7 97.4 63.7 39.3 42.5 147.9 92.6 22.4 22.4 36.8 196 678 513 0.8 1969 0.612 22.1 0.764 31.7 245.6 277.6 B 2 95.3 62 750 284 3175.972 77 01 30 20 106.1 -0.8 80.4 90.4 157.8 110.0 18.8 188 385 15.5 93.2 80.5 4.2 12.50 0.692 9.8 0.848 465 48.7 130.1 A 1 4.7 2.27 48 288 3181.807 77 02 07 11 155.7 155 56.7 392 1468 933 23.7 14.0 5 37.5 26.6 696 34.6 2.2 22.90 0.717 2.1 0.657 3.86 81.2 138.1 A P 1 7.6 3.21 350		205	19.7	36.84	0.716	0.369	1 000	284.89	603.03	6300
287 3181.807 77 02 07 11 1357 153 567 992 4468 933 254 105 375 20.6 968 946 22 225 0.671 21 0057 346 642 154 24 17 185 105 34 359 201 3199631 77 02 25 24 137.1 189 408 451 1591 1113 185 188 379 145 91.1 625 20 2.480 0.661 0.9 0.841 4.12 231.7 3560 FL 1 9.3 361 1100 300 344 152 231 1365 185 104 452 19 2235 0.680 26 0.715 3.76 2529 18.0 AR 1 75 2.82 160 0.4		207	17.9	24.45		0.484		1223.38	2589.55	3400
201 3262875 77 04 09 17 2173 139 631 649 1182 686 361 244 363 344 953 356 2.6 1.960 0.905 1.4 0.186 3.74 3154 19.1 A 1 6.7 2.72 44 0.6 204 3215894 77 04 18 2 1846 165 477 562 1633 1181 162 63 63 96 120 690 368 35 2.5 0.914 345 221.0 279 E 1 8.0 336 960 2.5		218	18.5	64.65		0.198	0.820		202.37	140
309 3286.869 77 05 23 24 310.2 397 29.6 29.8 92.4 55.0 499 499 39.8 48.5 97.2 72.3 06 5.209 0.886 87.5 1.012 9.41 181.4 61.7 AF 3 6.9 2.29 7 310 3299.809 77 06 05 15 228.2 47.9 12.8 15.7 160.7 11.25 17.0 10.5 36.3 13.0 76.1 36.8 2.6 2.070 0.516 18.4 1.02 3.14 195.7 74.1 BK 1 7.6 2.99 300 2.2 314 3341.837 77 01 7 34 2.409 - 13.6 67. 82.4 17.05 14.1 14.6 93 37.5 98 70.4 42.1 55 2.599 0.600 5.0 099 4.21 201.7 11.43 F 2 60 2.294 480 65		210	18.4	12.51	2.060	1.217		15813.00		87 000
314 3341.837 77 07 17 34 2405 -1.3 667 82.4 170.5 141.1 146 93 37.5 98 70.4 42.1 50.5 2299 70.68 91 30 0599 94 201.1 11-5 F 2 500 249 9400 50 317 3366.834 77 08 11 R 45.6 57.5 41.8 42.0 66.5 40.1 61.8 61.8 42.6 60.6 100.6 84.3 50 - 10.68 114.1 0.953 - 152.3 138.1 P A 1 6.1 1.72 1 325 3397.74 77 09 11 18 42.5 45.7 958.2 51.6 306 46.3 46.3 46.3 46.3 70.1 10.3 82.7 0.6 7.117 0.914 141.9 0.612 13.62 259.7 168.0 A 1 8.6 2.80 14		213	27.1	18.33	1.809	0.493		66757.19	141305.66	
331 3441882 77 10 25 6 1666 185 61.8 820 1699 396 133 70 250 69 714 311. 71 0.737 0.338 3.2 0.501 1.01 18.5 211.7 EL 2 91 4.13 10000 2500 133 1440749 77 11 01 14 10.08 494 442 445 714 414 882 582 398 595 114 0 90.1 04 4.275 0.800 1174 0472 808 2765 2203 A 1 7.7 253 9		225	21.1						68.90	
356 3455.787 77 11 08 0 151.2 29.8 79.5 79.8 31.1 18.0 69.8 69.8 41.2 68.7 110.5 83.2 2.1 9.328 0.898 150.4 0.952 17.70 156.7 225.4 G 1 9.1 3.60 73 339 3458.756 77 11 11 24 112.1 32.6 57.6 57.9 33.6 32.0 66.0 66.0 42.4 64.8 114.1 87.5 0.8 - 1.002 154.7 0.452 - 277.2 225.4 A 1 9.0 3.14 28				75.92		0.314				140
340 3458.739 77 11 11 R 57,1 14.2 40.6 42.0 134.5 81.1 29.1 25.6 37,1 26.8 88.1 69.1 0.9 2.126 0.810 59 0.465 3.85 10.87 48.4 8T EG 1 7.9 2.48 35 346 3467.09 77 11 22 17 16.7 48.7 4.3 63 16.25 10.82 15.7 8.5 35.2 11.1 64.1 31.0 2.3 1.60 0.443 11.5 0.892 2.31 22.6.8 239.4 E 3 8.6 35.0 150 160 349 342.290 77 12 0 5 R 12.3 4 0.453 457 72.6 42.2 58.0 58.0 40.1 57.0 0.46 66.8 0.8 4.585 0.592 12.84 c018 859 12.67 8.28 73.0 H 2 100 3.51 86		229	12.3	43.73	4.564	6.033		704.43	1491.07	150
379 3725807 78 08 05 R 369 552 429 432 708 426 600 600 422 588 979 818 0.4 - 1.033 1098 0.952 - 151.6 1321 P A 2 10.4 298 24 300 177881 78 08 08 R 79 574 391 391 698 410 586 386 402 574 981 834 0.3 6.489 0.831 10.6 0.966 12.01 1538 153 153 P A 1 72 2.17 4		231	27.9	52.72		0.352		1716.79	3633.96	4300
382 3770.852 78 08 10 R 43.5 562 32.5 32.6 66.3 39.2 60.2 60.2 40.8 591 39.0 75.9 0.3 10.088 0.906 114.3 0.952 19.22 150.7 157.0 P EF 3 9.5 2.74 14 383 3711.896 78 08 11 0 376 580 258 258 687 410.600 600 414 589 89.8 81.7 02 25800 0.962 111.5 0.981 50.51 159.1 158.0 P EF 1 9.0 228 5		232	35.0	434.10		0.000	0.447		4.13	130
384 3734.896 78 08 14 23 317.9 -23.5 76.1 83.3 150.8 98.2 21.1 12.1 36.6 18.1 80.3 53.8 5.8 2.166 0.677 4.0 0.699 3.63 76.5 320.8 A 1 6.8 2.99 280 0.8		235	19.1	171.01	14.343	7.863	0.887	4.08	8.64	720

Halliday et al., 1996

Gritsevich, 2009

50.27 0.878 0.630 0.826 198.96

241

16.7

421.15

730

Resolution

Adjusting the (vi, yi) values of Eq. [2] to the trajectory observed (vi, yi) values by means of a weighted least-squares method. Assigning manually the weighted factors may be quite complicated, so, since the height and velocity of a meteor decrease while it gets closer to the surface, the solution was proved to perform better if we take an exponential form of eq. [2] (see Gritsevich, 2008 for further details):

$$2\alpha \exp(-y) - \Delta \exp(-\beta) = 0, \quad \Delta = \overline{Ei}(\beta) - \overline{Ei}(\beta v^2)$$
$$\overline{Ei}(x) = \int_{-\infty}^{x} \frac{e^z dz}{z}$$

We can derive these new variables (α , β) for each meteoroid by minimizing this expression.

$$Q(\alpha, \beta) = \sum_{i=1}^{n} (F_i(y_i, v_i, \alpha, \beta))^2$$
$$F_i(y_i, v_i, \alpha, \beta) = 2\alpha \exp(-y_i) - \Delta_i \exp(-\beta) = 0$$

From Eq.[3], at the point were m=0 we have:

$$h_I = h_0 y_t = h_0 \ln 2\alpha\beta \quad [5]$$

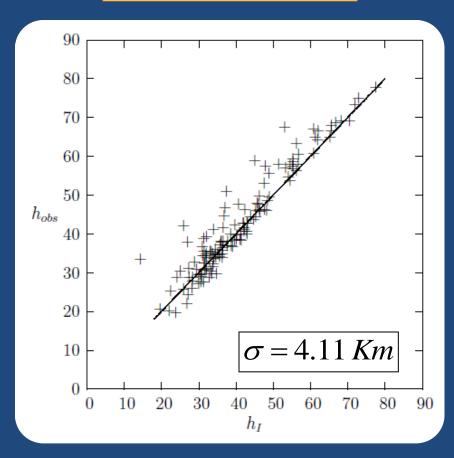
If we reorder Eq.[4] we have:

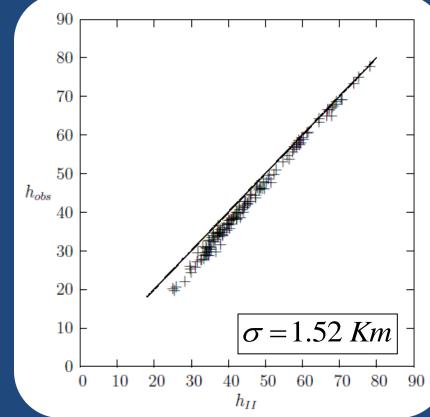
$$h_{II} = h_0 y_t = h_0 \ln\left(\frac{2\alpha\beta}{1 - e^{\beta(v_t^2 - 1)}}\right)$$
 [6]

Results

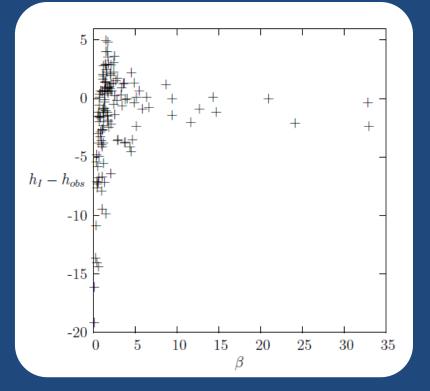
$$h_I = h_0 y_t = h_0 \ln 2\alpha\beta$$

$$h_{II} = h_0 y_t = h_0 \ln\left(\frac{2\alpha\beta}{1 - e^{\beta(v_t^2 - 1)}}\right)$$



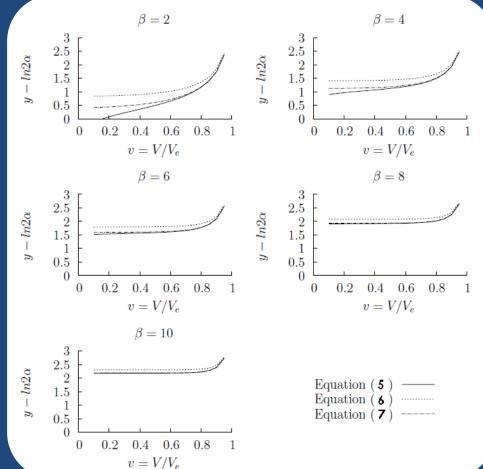


Results (II)



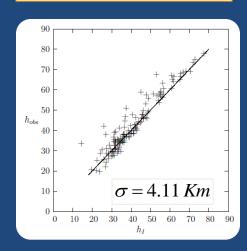
Differences could be due to the use of eq. [4] established for high β .

As suggested by Gritsevich et al. (2015), we would rather use the modification: $\beta \rightarrow \beta - 1.1$

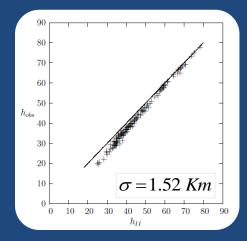


$$h_{III} = h_0 y_t = h_0 \ln\left(\frac{2\alpha(\beta - 1.1)}{1 - e^{(\beta - 1.1)(v_t^2 - 1)}}\right) \quad [7]$$

$$h_I = h_0 y_t = h_0 \ln 2\alpha\beta$$



$$h_{II} = h_0 y_t = h_0 \ln\left(\frac{2\alpha\beta}{1 - e^{\beta(v_t^2 - 1)}}\right)$$



$$h_{III} = h_0 y_t = h_0 \ln \left(\frac{2\alpha(\beta - 1.1)}{1 - e^{(\beta - 1.1)(v_t^2 - 1)}} \right)$$

Discussion

For $\beta > 5$ eq. [5] shall give good results. But the decrease in v near the terminal point is not considered.

Eq.[6] shows a lineal tendency. Still discrepancies in all β values.

The modification made in [7] leads to a good agreement between observed and theoretical data. However at low β , some discrepancies appear.

The inverse problem is possible for non decelerated bodies if the terminal height is known \rightarrow constraints in α and β .

Meteor height as a function of time → new problems may be scoped:
Determination of luminous efficiency based on meteor duration.
Critical Kinetic Energy to produce luminosity.

The small discrepancies at low β shall be taken into account for future planetary defense applications. Meteoroids could reach lower height than those predicted. We use dimensionless parameters instead of the empirical set A,B, C coefficients of Ceplecha and McCrosky (1976). However α and β keep the same variable dependency as the PE criterium. • We have derived the terminal heights for MORP fireballs using a new developed methodology.

• This methodology had only been tested on several fully ablated fireballs with large β values (Gritsevich and Popelnskaya, 2008).

• We were particularly interested in determining whether this new mathematical approach works equally for fully ablated fireballs and meteorite-producing ones.

• We introduced a new modification in the methodology which allows to get a higher accuracy.

• We foresee a calculation of terminal height to be useful when the lower part of the trajectory was not instrumentally registered.

• It also brings critical knowledge into the problem when one needs to predict how long will be a total duration of the luminous flight or at which height a fireball produced by a meteoroid with given properties would terminate.

• Based on our investigations we can highly recommend the use of equation [7] also to solve inverse problem when terminal height and velocity are available from the observations, and parameters α and β need to be derived.



M. Moreno-Ibáñez, et al., Icarus 250, 544-552 (2015).

mmoreno@ice.csic.es, trigo@ice.csic.es, maria.gritsevich@nls.fi

Z. Ceplecha and R.E. McCrosky, J.Geophys. Res. 81, 6257-6275 (1976).
M. Gritsevich, Solar Syst. Res. 41 (6), 509-514 (2007).
M. Gritsevich, Solar Syst. Res. 42, 372-390 (2008).
M. Gritsevich and N.V. Popelenskaya, Doklady Phys. 53, 88-92 (2008).
M. Gritsevich, Adv. Space Res. 44. 323-334 (2009).
M. Gritsevich et al., Mat. Model. 27 (2) 25-33 (2015).
I. Halliday et al., Meteorit. Planet. Sci. 31, 185-217 (1996).
V.P. Stulov et al., Aerodinamika bolidov, Nauka (1995).
V.P. Stulov, Appl. Mech. Rev. 50, 671-688 (1997).
V.P. Stulov, Planet. Space Sci. 46, 253-260 (1998).
V.P. Stulov, Planet. Space Sci. 52 (56), 459-463 (2004).
G.W. Wetherill and D.O. Revelle, Icarus 48, 308-328 (1981).